

SELF-CONSISTENT LCAO LOCAL DENSITY DETERMINATION OF ANISOTROPIC COMPTON PROFILE AND X-RAY STRUCTURE FACTORS IN DIAMOND*

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Our newly developed fully Self-Consistent Numerical Discrete Variational Method (which exactly treats all non-spherical parts of the crystal potential) yields X-ray scattering factors and anisotropies in the Compton profile of diamond in very good agreement with experiment.

Considerable interest exists in the predictions of the local density formalism (LDF) [1] on the ground state properties of solids, e.g. cohesive energies [2], X-ray form factors [3, 4] and Compton profiles [5]. LDF based calculations are usually beset with the difficulties of solving self-consistently the associated one-particle equation characterized by a *multi-center* non-spherical potential and hence, a variety of approximations have been introduced to reduce the complexity of the problem, LCAO-type calculations [4, 5] have overcome the difficulty of treating non-muffin-tin potentials and have demonstrated that efficient convergence with respect to the size of the basis set [4] can be obtained. However, the problem of carrying this type of calculation or self-consistency (SC) still remains a formidable task. Although these methods are capable of yielding reasonable results for the eigenvalues, an accurate evaluation of ground state functionals of the electron density is still non-trivial.

In this paper we apply our newly developed [6] self-consistent numerical discrete variational method (DVM) [7] to study the X-ray scattering factors and directional Compton profile for diamond. All non-spherical parts of the crystal potential are treated exactly using an efficient numerical LCAO basis set and a numerical Diophantine integration scheme. Local density exchange and correlation are incorporated directly into the crystal potential and full self-consistency is obtained. The resulting X-ray form factors are in good agreement with experiment [8, 9]

and with previous Hartree-Fock calculations [10]. While the non-SC Compton profile is too high at low momenta and lacks some high momentum components, the fully self-consistent results agree very well with experiment. Unlike the Hartree-Fock results, the anisotropy of the profile is found to be in reasonable agreement with experiment.

For the crystal problem, the general potential is given by

$$V(\mathbf{r}) = V_{\text{coul}}(\mathbf{r}) + F_{\text{ex}}[\rho_{\text{sup}}(\mathbf{r})] + F_{\text{corr}}[\rho_{\text{sup}}(\mathbf{r})] \quad (1)$$

with the exchange, F_{ex} , and correlation, F_{corr} , potentials given in terms of the local density functions written in terms of superposed overlapping atomic densities, $\rho_{\text{sup}}(\mathbf{r})$. We use the free-electron $\rho^{1/3}$ exchange potential for F_{ex} and the correlation energy functional of Singwi et al. [11] as fitted to analytic form by Hedin and Lundqvist [12]. We do not spherically average $V(\mathbf{r})$ or linearize the local density functionals. The crystal wave functions $\psi_j(\mathbf{K}, \mathbf{r})$ are expanded in terms of Bloch functions $\Phi_{\mu}^{\alpha}(\mathbf{K}, \mathbf{r})$ [which are given in terms of LCAO basis orbitals χ_{μ}^{α}]. Unlike previous efforts which used simple analytic basis functions to overcome difficulties in calculating many-center integrals appearing in matrix elements (ME) of $V(\mathbf{r})$, we are able to exploit the variational efficiency of accurate *numerical* basis functions because we do not employ any analytic algorithms for calculating ME. Thus our $\chi_{\mu}^{\alpha}(\mathbf{r})$ are determined as numerical solutions of the atomic potential equivalent of eq. (1). Details of the method, and evidence for the variational superiority of even a minimal set (e.g., 1s, 2s and 2p) to a double-zeta Slater basis are given elsewhere

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